



Introduction to complex fluids

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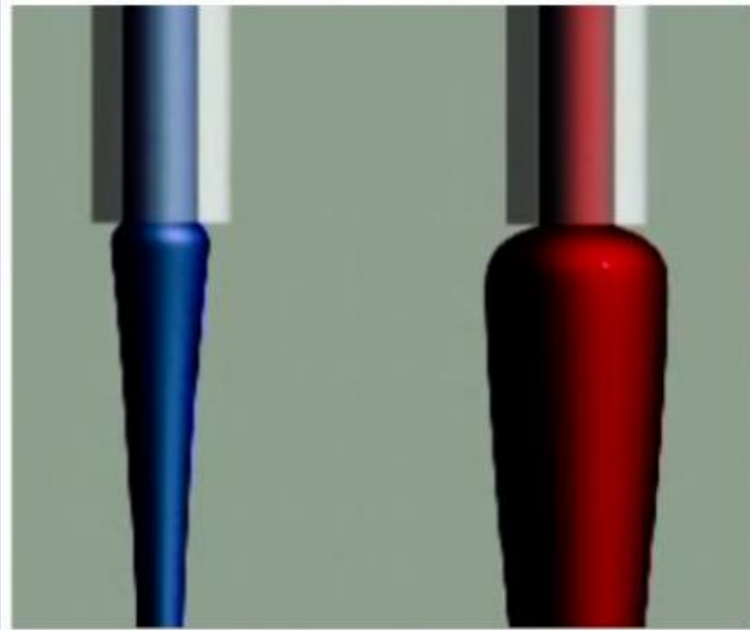
Overview

- Newtonian vs Non-Newtonian (complex) fluids
 - Complex rheological models
 - Generalized Newtonian fluids
 - Viscoelastic fluids
 - Chemorheological fluids (Polyurethane foams)
- The Navier Stokes Equation
- The Navier-Stokes Brinkman Equations

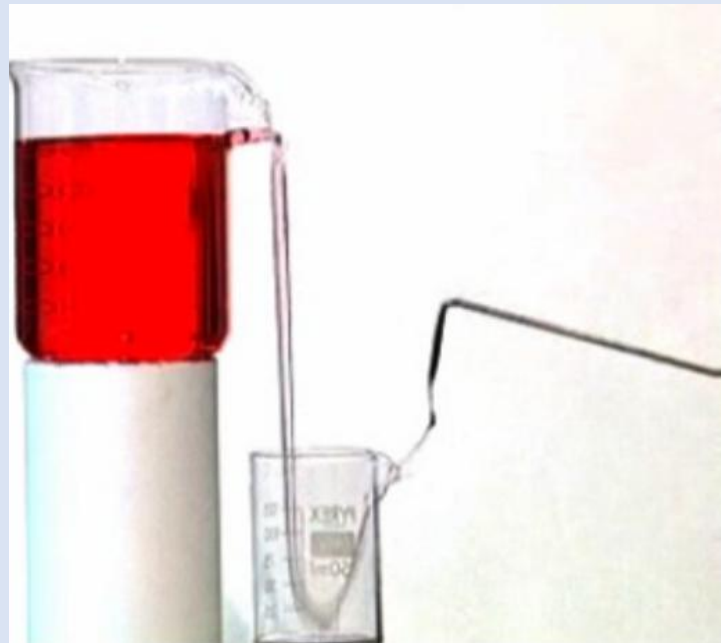
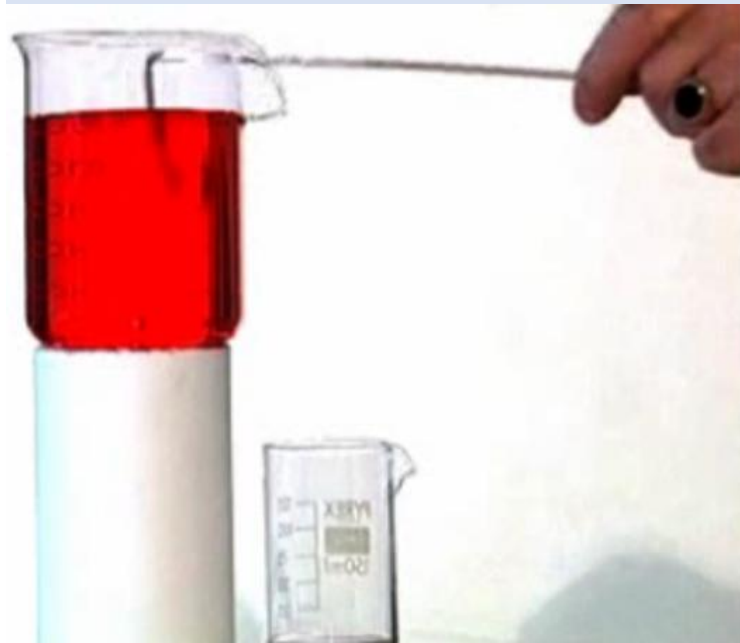
Simple Fluids



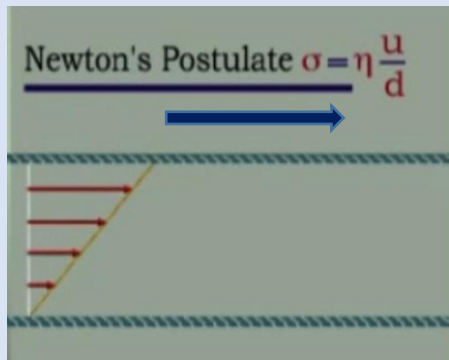
Simple vs Complex Fluids



Complex Fluids



Newtonian Fluids



σ is the shear stress

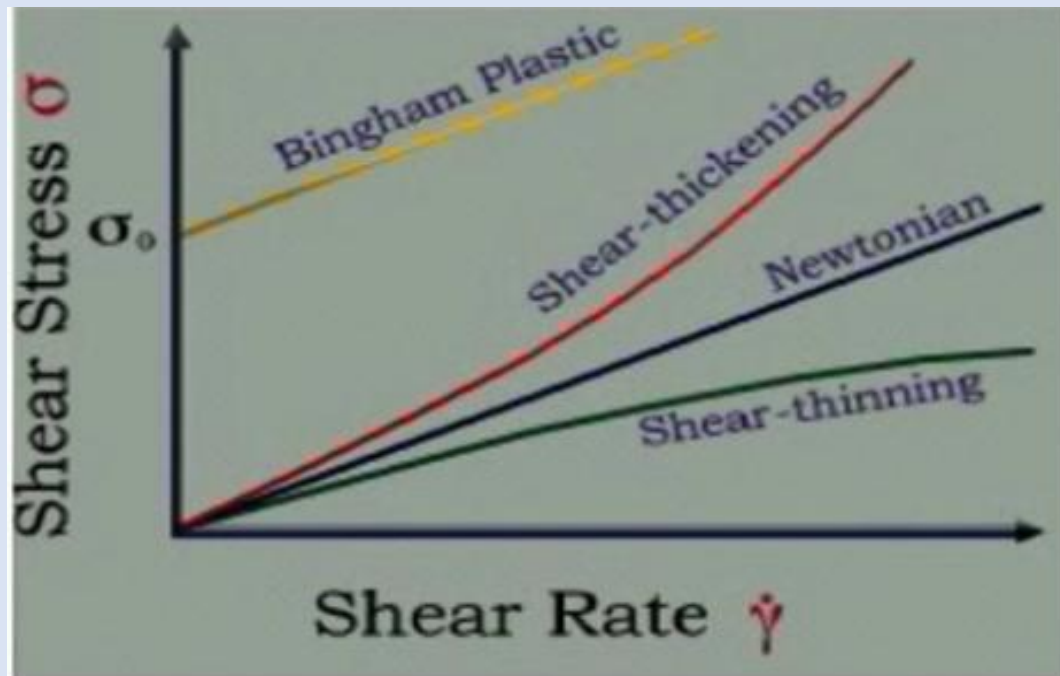
$\frac{u}{d}$ is the velocity gradient

η is the coefficient of viscosity

Table 1.1: Temperature dependent viscosity models

Reynold's model	$\eta(T) = \eta_0 \exp(-BT)$
Vogel's model	$\eta(T) = \eta_0 \exp\left(\frac{B}{T-T_0}\right)$
Arrhenius model	$\eta(T) = \eta_0 \exp(E/RT)$
Nahme model	$\eta(T) = \eta_0 \exp\left(\frac{T-T_0}{T_0}\right)$
William-Landel-Ferry (WLF) model	$\eta(T) = \eta_0 \exp\left(\frac{-C_1(T-T_0)}{C_2+T-T_0}\right)$
Fulcher model	$\log_{10}(\eta(T)) = -A + \frac{B \times 10^3}{T-T_0}$

Stress strain Curves: Simple vs Complex Fluids



- Viscoelasticity
- Variable Viscosity
- Normal Stresses
- Extensional Viscosity
- Thixotropy and Anti-Thixotropy
- Other Effects

The Navier-Stokes Equations

The **Navier-Stokes Equations** subject to appropriate boundary conditions

- **Conservation of mass** $\nabla \cdot \vec{v} = -\frac{1}{\rho} \left(\frac{\partial \rho}{\partial t} + \vec{v} \cdot \nabla \rho \right)$
- **Conservation of linear momentum** $\frac{\partial \rho \vec{v}}{\partial t} + \nabla \cdot (\rho \vec{v} \vec{v}) = -\nabla p + \nabla \cdot \mathbf{T} + \mathbf{F}$
- **Conservation of energy:** $\rho C_p \left(\frac{\partial T}{\partial t} + \vec{v} \cdot \nabla T \right) = \nabla \cdot (\kappa \nabla T) + \frac{1}{2} (\mu_{mix} \mathbf{D} : \mathbf{D})$

Constitutive relations

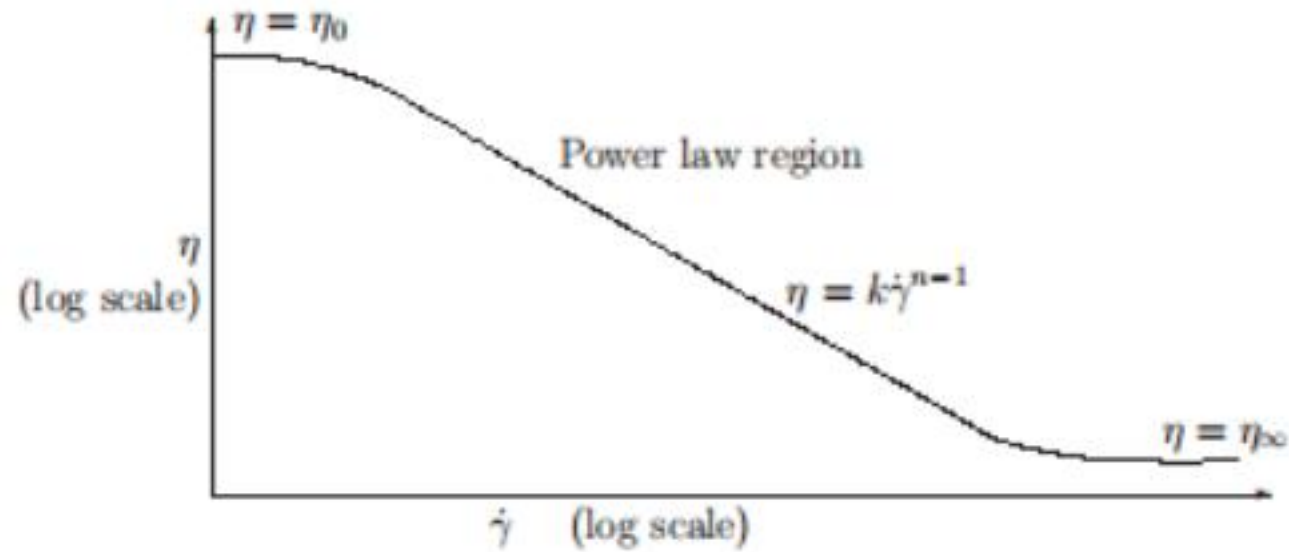
- Newtonian $\mathbf{T} = 2\eta\mathbf{D}$, with, $\mathbf{D} = \frac{1}{2} [\nabla\mathbf{v} + (\nabla\mathbf{v})^T]$.

- Generalized Newtonian** $\mathbf{T} = 2\eta(\dot{\gamma}) \mathbf{D}$.

- Power law Model $\eta = K\dot{\gamma}^{n-1}$,
 - if $n < 1$ $\lim_{\dot{\gamma} \rightarrow 0} \eta(\dot{\gamma}) = \infty$ $\lim_{\dot{\gamma} \rightarrow \infty} \eta(\dot{\gamma}) = 0$,
 - if $n > 1$ $\lim_{\dot{\gamma} \rightarrow 0} \eta(\dot{\gamma}) = 0$ $\lim_{\dot{\gamma} \rightarrow \infty} \eta(\dot{\gamma}) = \infty$.

- Prandtl-Eyring $\eta = \eta_0 \frac{\sinh^{-1}(\lambda\dot{\gamma})}{\lambda\dot{\gamma}}$, $\lim_{\dot{\gamma} \rightarrow 0} \eta(\dot{\gamma}) = \eta_0$, $\lim_{\dot{\gamma} \rightarrow \infty} \eta(\dot{\gamma}) = 0$.

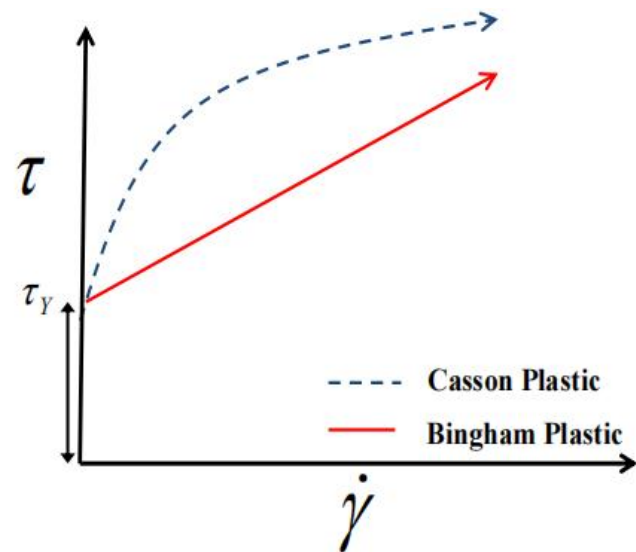
Viscosity shear rate relationship: Shear thinning fluid



More generalized Newtonian fluids

- Powell-Eyring model $\eta = \eta_\infty + (\eta_0 - \eta_\infty) \frac{\sinh^{-1}(\lambda\dot{\gamma})}{\lambda\dot{\gamma}}$. $\lim_{\dot{\gamma} \rightarrow 0} \eta = \eta_0$ and $\lim_{\dot{\gamma} \rightarrow \infty} \eta = \eta_\infty$.
- Cross model $\frac{\eta - \eta_\infty}{\eta_0 - \eta_\infty} = \frac{1}{1 + (\lambda\dot{\gamma})^m}$, $\frac{\eta_0 - \eta}{\eta - \eta_\infty} = (\lambda\dot{\gamma})^m$.
 if $m < 1$ $\lim_{\dot{\gamma} \rightarrow 0} \eta(\dot{\gamma}) = \eta_0$ $\lim_{\dot{\gamma} \rightarrow \infty} \eta(\dot{\gamma}) = \eta_\infty$,
 if $m > 1$ $\lim_{\dot{\gamma} \rightarrow 0} \eta(\dot{\gamma}) = \eta_\infty$ $\lim_{\dot{\gamma} \rightarrow \infty} \eta(\dot{\gamma}) = \eta_0$,
- Sisko model $\eta = \eta_\infty + K\dot{\gamma}^{n-1}$.
- Carreau-Yasuda model $\frac{\eta - \eta_\infty}{\eta_0 - \eta_\infty} = [1 + (\lambda'\dot{\gamma})^c]^{(n-1)/c}$. $c = 2$ (Carreau Model)

Yield Stress Fluids



Bingham plastic:

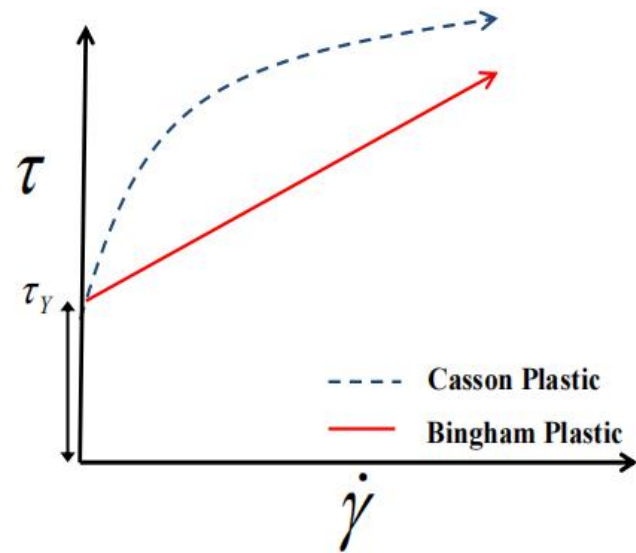
$$\tau = \tau_Y + \eta_p \dot{\gamma}.$$

$$\eta = \begin{cases} \infty & \text{if } \tau \leq \tau_Y, \\ \eta_0 + \frac{\tau_Y}{\dot{\gamma}} & \text{if } \tau > \tau_Y. \end{cases}$$

Herschel-Bulkley model:

$$\eta = \begin{cases} \infty & \text{if } \tau \leq \tau_Y, \\ K \dot{\gamma}^{n-1} + \frac{\tau_Y}{\dot{\gamma}} & \text{if } \tau > \tau_Y. \end{cases}$$

Yield Stress Fluids



Casson plastic:

$$\sqrt{\eta} = \begin{cases} \infty & \text{if } \tau \leq \tau_Y, \\ \sqrt{\eta_0} + \sqrt{\frac{\tau_Y}{\dot{\gamma}}} & \text{if } \tau > \tau_Y. \end{cases}$$

Viscoelastic fluids

- Jeffery model
- Maxwell Fluid
- Order grade fluid
- Differential Viscoelastic Model

$\sigma_{12} = \dot{\gamma}\eta(\dot{\gamma}),$ Shear Stress,
 $\sigma_{11} - \sigma_{22} = N_1(\dot{\gamma}),$ First Normal stress difference,
 $\sigma_{22} - \sigma_{33} = N_2(\dot{\gamma}),$ Second Normal stress difference.

Table 1.2: Some differential viscoelastic models

Model Name	Constitutive model for σ	N_1 and N_2
Oldroyd-B	$\sigma + \lambda \overset{\nabla}{\sigma} = 2\eta_v \mathbf{D}$	$N_1 > 0$ $N_2 \equiv 0$
Geisekus	$\sigma + \alpha (\sigma \cdot \sigma) + \lambda \overset{\nabla}{\sigma} = 2\eta_v \mathbf{D}$	$N_1 > 0$ $N_2 \neq 0$
Johnson-Segalman	$\sigma + \lambda \overset{\square}{\sigma} = 2\eta_v \mathbf{D}$	$N_1 > 0$ $N_2 \neq 0$
Diffusive Johnson-Segalman	$\sigma + \lambda \overset{\square}{\sigma} = 2\eta_v \mathbf{D} + \varepsilon \nabla^2 \sigma$	$N_1 > 0$ $N_2 \neq 0$
	$\sigma + \lambda \overset{\square}{\sigma} = 2\eta_v \mathbf{D} + \varepsilon \nabla^2 \mathbf{D}$	$N_1 > 0$ $N_2 \neq 0$
Phan-Thien-Tannar	$\exp\left(\pi \frac{\lambda}{\eta_v} \text{tr}(\sigma)\right) + \sigma + \lambda \overset{\square}{\sigma} = 2\eta_v \mathbf{D}$	$N_1 > 0$ $N_2 \neq 0$

Chemorheological fluids

$$\eta = \eta(\theta, p, \dot{\gamma}, t, F, \alpha)$$

$$\mathbf{T} = \eta_m \mathbf{D}$$

$$\eta_m = \begin{cases} \eta_F, & \text{for } \varphi \geq \varphi_s, \\ \eta_F \frac{\varphi}{\varphi_s} + \eta_g \left(1 - \frac{\varphi}{\varphi_s}\right), & \text{for } 0 \leq \varphi < \varphi_s. \end{cases}$$

$$\eta_F = \eta_{oo} \exp\left(\frac{E_\eta}{RT}\right) \cdot \left(\frac{\zeta_c}{\zeta_c - \zeta}\right)^{h(\zeta)} f(\varphi_g),$$

THANK YOU
FOR YOUR ATTENTION